# RESERVOIR AND PRODUCTION ENGINEERING OF WELL A-18 IN THE LOS AZUFRES GEOTHERMAL FIELD IN MEXICO 

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## ABSTRACT

Reservoir and production engineering tools are applied to data obtained in a pressure buildup test on well A-18 in the Los Azufres geothermal field in México. The well produces a twophase mixture through a single fracture in a double-porosity media. The reservoir parameters were estimated as follows: permeability-thickness between 5.4 to $8.1 \times 10^{-12} \mathrm{~m}^{3}(18,000$ to $27,000 \mathrm{md}-\mathrm{ft})$; the ratio of fissure system storativity to that of the total fissure-matrix system 0.1. Test interpretation showed the presence of a sealing boundary near the well. Using estimated reservoir parameters the simulation of an idealized well producing at constant mass flow rate near a linear boundary was done considering both homogeneous and double-porosity reservoirs.

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## 1. INTRODUCTION

At Well A-18 in the Los Azufres Geothermal Field, there will be installed a $5 \mathrm{MW}_{\mathrm{e}}$ non-condensing turbo-generator unit, to be commissioned in 1987, according to México's National Geothermoelectric Expansion Program.

In order to know the initial characteristic of this well before starting its exploitation, several tests have been carried out: pressure transient tests, and production tests.

The importance of these and other tests is because their interpretation can be used to predict the well's behaviour and also permits the taking of quick decisions to repair a well or drill another one.

The above mentioned is applied directly to a specific well, however, the reservoir parameters obtained from the tests must be utilized in the total reservoir model to predict the complete system behaviour and also to decide between different alternatives.

In the present work, practical tools to estimate reservoir parameters from a pressure buildup test are presented.

## 2. LOS AZUFRES FIELD

Geothermal prospecting studies in México started at the beginning of the 1950s Alonso (1985). At the present, there are more than 60 known areas which discharge geothermal fluids on the surface. On Figure 1 are presented some of these, taking into account their importance.

The Pathe zone in the state of Hidalgo was the first geothermal area where wells were drilled to generate electricity. This was done because of the relatively short distance to México City, where the consumption of electrical energy has always been high. The mass flow rate was poor and the heat content (specific enthalpy) low. However, there was installed the first geothermoelectric plant, generating about 600 kW . The importance of Pathé is because it was demonstrated to be feasible to exploit the geothermal resource to generate electricity.

The four most important fields in México are: Cerro Prieto, Los Azufres, Los Humeros and La Primavera. Razo (1985) presents the geological, geophysical and geochemical characteristics of these and other zones. Molinar (1985) shows a general view about the evaluation of these fields.

Well A-18, was drilled in the Los Azufres geothermal field, which is located in the state of Michoacán (Reyes, 1985), 200 km northwest from México City. The reservoir is classified as liquiddominated and has a surface area of about $30 \mathrm{~km}^{2}$. More than 50 wells have been drilled (Figure 2). Although the studies indicate the existence of one reservoir, the field has been divided into two parts; namely, north and south zones. In the north zone, most of the producing wells discharge a two-phase mixture. On the other hand, in the south zone there are about 4 wells which produce superheated steam. The specific enthalpy ranges from 1000 to $2850 \mathrm{~kJ} / \mathrm{kg}$, and the best steam well produces about $30 \mathrm{~kg} / \mathrm{s}$.

The reservoir has been commercially exploited since 1982 , using

5 non-condensing turbo-generator units (Ortega, 1985). The nominal capacity of each unit is 5 MWe and the admission steam pressure and mass flow rate are 0.8 MPa and $16.4 \mathrm{~kg} / \mathrm{s}$, respectively. The separated water is injected into the subsurface to avoid ecological deterioration.

Alonso (1985) has presented the expected National Geothermoelectric Expansion Program. For the case of Los Azufres the installation of some 7 small-scale turbo-generator units is planned, similar to those already installed, and also the installation of the large-scale Tejamaniles Geothermoelectric Central in the south zone. The nominal capacity of the central is 50 MW . According to this program, and considering the flow characteristics of all the wells, it was decided to use the well A-18 to supply one of this small units.

## 3. WELL CHARACTERISTICS

### 3.1. Well Completion

Well A-18 was drilled in the south zone of the Los Azufres field, close to an area where most of the wells produce superheated steam at relatively shallow depths (about 700 m ). It was completed at 1328 m depth using the following pipes: $13-3 / 8^{\prime \prime}$ from 0 to $300 \mathrm{~m} ; ~ 9-5 / 8^{\prime \prime}$ from 0 to $1000 \mathrm{~m} ; 7^{\prime \prime}$ from 959 to 1328 m . In Figure 3 is presented the completion of this well, the beginning of the slotted can be easily distinguished at 1013 m depth. Completions like this are common at Los Azufres.

### 3.2. Output Curve

In March 1986 the output characteristic curve of this well was obtained. The maximum total mass flow rate was calculated to be as high as $43 \mathrm{~kg} / \mathrm{s}$; the specific enthalpy of the mixture as $1764 \mathrm{~kJ} / \mathrm{kg}$; while the wellhead pressure was measured 0.9 MPa (Table 1). In Figure 4 are presented the steam and water characteristic curves, both at atmospheric conditions. In Table 1 the specific enthalpy of the mixture increases as the mass flow rate increases too. This effect occurs when there is a heat transfer process from the rock matrix to the fluid. In Table 1 and Figure 4 is possible to appreciate that the production of the well is controlled by the reservoir.

## 4. PRESSURE BUILDUP MEASUREMENTS

### 4.1. Test Description

On March 17, 1986, well A-18 was opened to carry out a pressure buildup test. According to field experiences acquired for several years in this matter, the well was left to produce for two days in order to reach a stable production state before the test was commenced. The flow rate was controlled by using a 2 inch orifice, installed close to the wellhead in the pipeline to the silencer. Since the production started, the common surface measurements were periodically taken and registered. Few hours before the test started, pressure and temperature logs were run to know the state of the fluid flowing into the well. After that, the test was developed using both temperature and pressure recorders.

### 4.2. Measurements Before Shut-In

During the drawdown period (production), three basic parameters were registered at the surface: the wellhead pressure $P_{w n}$, critical lip pressure $P_{c}$ and head of water in weir box. The measured quantities were

$$
\begin{aligned}
\mathrm{P}_{\mathrm{w}} & =3.2 \mathrm{MPa} \\
\mathrm{P}_{\mathrm{c}} & =0.077 \mathrm{MPa} \\
\mathrm{~A} & =0.138 \mathrm{~m}
\end{aligned}
$$

The diameter of the discharge pipe was 0.1985 m .

Simultaneous pressure and temperature logs were run downhole. The measured pressure and temperature values are reported in Table 2 and their plots against depth are shown in Figures 5 and 6, respectively. These data were also used to construct Figures 7 and 8 which will be discussed later.

### 4.3. Measurements During Buildup

Before the well was shut-in, pressure and temperature elements were lowered to 1200 m depth. Kuster elements recorded the pressure and temperature behaviour during a 19 hours test period. After that, the elements were brought to the surface, and the registered deflections on the metallic charts were read. The information obtained in this test is presented in Table 3. Pressure and temperature data were also plotted against the elapsed time and they are shown in Figures 9 and 10.

## 5. FLOWING CONDITIONS

### 5.1. Specific Enthalpy and Mass Flow Rate

To estimate the mass flow rate and the specific enthalpy of the fluid during the discharging period (drawdown), the empirical relation of James (1962) was employed. That expression is written in S.I. units as

$$
\begin{equation*}
w=\operatorname{Exp}(16.394) D_{d}^{2} P_{c} 0.96 h_{0}^{-1.102} \tag{1}
\end{equation*}
$$

For the case of Los Azufres, where the atmospheric pressure is 0.073 MPa , one has from steam tables (Keenan et al., 1978) the following

$$
\begin{aligned}
\mathrm{v}_{1} & =1.0368 \mathrm{E}-03 \mathrm{~m}^{3} / \mathrm{kg} \\
\mathrm{~h}_{1 \mathrm{~s}} & =2280.48 \mathrm{~kJ} / \mathrm{kg} \\
\mathrm{~h}_{\mathrm{s}} & =2661.8 \mathrm{~kJ} / \mathrm{kg}
\end{aligned}
$$

where, $v_{1}$ is the specific volume of liquid water, $h_{1} s$ is the latent heat content, and $h_{s}$ the steam heat content. To determine the liquid water flow rate through a $V$-notch ( $90^{\circ}$ ) weir box (ASME, 1971), one can use

$$
\begin{equation*}
\mathrm{w}_{1}=1.3345 \mathrm{~A}^{2} .475 / \mathrm{V}_{1} \tag{2}
\end{equation*}
$$

where $A$ is the head of water in $m$, and $w_{1}$ the water flow rate in $\mathrm{kg} / \mathrm{s}$. Substituting $\mathrm{v}_{1}$ into Equation 2, one gets

$$
\begin{equation*}
\mathrm{w}_{1}=1287.16 \mathrm{~A}^{2} .475 \tag{3}
\end{equation*}
$$

The steam fraction is defined at atmospheric conditions as

$$
\begin{equation*}
x=w_{s} /\left(w_{1}+w_{s}\right)=\left(h_{0}-h_{1}\right) / h_{1} s \tag{4}
\end{equation*}
$$

where, $W_{s}$ is the steam flow rate in $\mathrm{kg} / \mathrm{s}$, $\mathrm{h}_{0}$ is the specific enthalpy of the steam-water mixture in $\mathrm{kJ} / \mathrm{kg}$ (assuming that the stagnation and the steam-water mixture enthalpies are equal),
and $h_{1}$ the liquid water heat content in $\mathrm{kJ} / \mathrm{kg}$. From Equation 4 one has

$$
\begin{equation*}
w=w_{1} h_{1 s} /\left(h_{s}-h_{0}\right) \tag{5}
\end{equation*}
$$

Substituting Equation 3 into Equation 5 with the values of the enthalpies for Los Azufres one gets

$$
\begin{equation*}
w=\operatorname{Exp}(14.8923) A^{2.475} /\left(2661.8-h_{0}\right) \tag{6}
\end{equation*}
$$

The final expression in terms of the stagnation enthalpy is derived from Equation 1 and 6, and can be written as

$$
\begin{equation*}
F\left(h_{0}\right)=a_{1} D_{d}^{2} P_{c} 0.96 h_{0}^{-1.102}\left(2661.8-h_{0}\right)-A^{2.475}=0 \tag{7}
\end{equation*}
$$

where

$$
a_{1}=4.487
$$

The latter equation must be solved for $h_{0}$ using a mathematical convergence method (e.g. Newton-Raphson). After that process, one has to come back to the expression which contains the mass flow rate. The computed values for both; the specific stagnation enthalpy; and the mass flow rate were found, respectively

$$
\begin{aligned}
& \mathrm{h}_{\mathrm{o}}=1314 \\
& \mathrm{w}=16.2 \mathrm{~kJ} / \mathrm{kg} \\
& \mathrm{~kg} / \mathrm{s}
\end{aligned}
$$

The lip pressure was not corrected for the presence of gas in the mixture because the amount of gas present was not available but is it generally low in Los Azufres.

### 5.2. Pressure and Temperature Logs

Returning to the flowing pressure profile in Figure 5, it is possible to distinguish three different straight lines. The first straight line starts at the surface and finishes where the intersection with the second line occurs at 950 m depth. The change in slope between the first and the second straight
lines is due to the reduction in pipe diameter. The intersection of the second and third straight lines is located at 1225 m depth. Here the change in slope is related to the feed zone. Flowing temperature profile (Figure 6) shows an almost typical temperature distribution inside a geothermal well. However at about 1250 m depth there is a small temperature inversion which also can be related to the feed zone.

### 5.3. Fluid Thermodynamic State

In order to know the state of the fluid throughout the well during flowing, the Clapeyron Diagram in Figure 7 was utilized. Pressure and temperature data under flowing conditions were plotted on the diagram. It was possible to determine that twophase flow occurs between from approximately 1225 m depth and to the surface. The presence of a single liquid water phase was detected below that depth.

### 5.4. Wellbore Simulation

Taking into account the calculated specific enthalpy and mass flow rate as well as the measured wellhead pressure and the well completion, a program used by Ambastha and Gudmundsson (1986), was run from top to bottom. The data and output are presented in Appendix A. The measured and calculated pressures can be observed and compared in Figure 8. This plot shows that the simulator gave reasonable fit. Although not presented, the temperature profile was also a good fit.

### 5.5. Feed Zone Depth

After analyzing the information available and employing the results of the simulator the feed zone was determined to be between 1200 and 1250 m depth. Below that depth, the pressure increases according to a hydraulic column, that can be due to a small or non-production horizon.

## 6. INTERPRETATION OF BUILDUP TEST

### 6.1. Pressure and Temperature Buildup

In the Clapeyron Diagram (Figure 11) is presented the thermodynamic behaviour of the fluid at 1200 m depth during the buildup test. The thermodynamic state of the fluid was in the compressed liquid region.

### 6.2. Identification of Model

### 6.2.1. Inner boundary

Pressure transient test interpretation starts by plotting the pressure increment $\delta \mathrm{P}$ against the time increment $\delta \mathrm{t}$ a $\log -\log$ paper as shown in Figure 12. The inner boundary effect is felt at the earlier elapsed time (Gringarten, 1985). The dominant effect can be: wellbore storage, skin, fracture, and partial penetration.

Wellbore storage effect is due to expansion of the fluid inside the well and is characterized by a straight line of one unit slope in the diagnostic plot (Figure 12). Skin effect is due to the presence of some damage in the walls of the well and it produces a steady state pressure drop. Partial penetration results from uncompleted drilling process through the total reservoir thickness (normally found in geothermics). Fractures exhibit on a log-log plot a straight line with one-half slope when it has a very high conductivity or one-quarter slope when the conductivity is low.

The inner boundary was determined to be a single medium conductivity fracture because the slope of the fitted straight line which passes through the earliest points (Figure 12) is between the one-half and one-quarter slopes.

### 6.2.2. Reservoir behaviour

Gringarten (1985) defines reservoirs according to their response, as homogeneous or heterogeneous. These responses can be very similar when the data are not properly plotted on log-log paper (scale problem). However, it is possible to avoid that confusion by plotting the derivative $\mathrm{d}(\delta \mathrm{P}) / \mathrm{d}(\operatorname{Ln}(\delta t))$ against $\delta \mathrm{t}$ on the same kind of paper.

This idea was considered to apply directly to the test data, however, there was noise in the results and that made it difficult to distinguish clearly the reservoir response. Least squares method was used to smooth the data. In Figure 13 is presented 8 P against $\operatorname{Ln}(\delta t)$ on Cartesian axis for both cases: measured and fitted data. The fitted curve is

$$
\begin{equation*}
\delta P=A_{1}+A_{2} * Z+A_{3} * Z^{2}+A_{4} * Z^{3}+A_{5} * Z^{4} \tag{8}
\end{equation*}
$$

where

$$
\begin{aligned}
\mathrm{A}_{1} & =3.634032 \\
\mathrm{~A}_{2} & =-2.113564 \\
\mathrm{~A}_{3} & =0.4576055 \\
\mathrm{~A}_{4} & =-4.222444 \mathrm{E}-02 \\
\mathrm{~A}_{5} & =1.415709 \mathrm{E}-03 \\
\mathrm{Z} & =\operatorname{Ln}(\delta \mathrm{t})
\end{aligned}
$$

with
Coefficient of determination $=0.995$
Coefficient of correlation $=0.998$
Standard error of estimate $=2.79 \mathrm{D}-03$

Therefore the fit was reasonable, and derivative function becomes

$$
\begin{equation*}
\mathrm{d}(\delta \mathrm{P}) / \mathrm{d}(\operatorname{Ln}(\delta t))=\mathrm{A}_{2}+2 * \mathrm{~A}_{3} * \mathrm{Z}+3 * \mathrm{~A}_{4} * \mathrm{Z}^{2}+4 * \mathrm{~A}_{5} * \mathrm{Z}^{3} \tag{9}
\end{equation*}
$$

Plotting the derivative function as described before, Figure 14 was obtained, which shows the characteristic hump of a heterogeneous reservoir response (Gringarten 1985). Now going back to Figure 12, one can appreciate in the infinite acting
period two straight lines with the same slope, which is typical of double-porosity reservoirs.

### 6.2.3. Outer boundary

At late time, it is possible to observe a faster pressure increment in the diagnostic plot (Figure 12), which is the characteristic of an outer boundary. To determine the type of boundary, it was necessary to construct the specialized plot shown in Figure 15. One can see that the slope of the late time straight line is twice the total system reservoir slope. This is due to the existence of a sealing fault near the well.

### 6.2.4. Complete reservoir behaviour

The complete behaviour is obtained by combining the individual behaviors. In that way, the complete behaviour is defined as: a single medium conductivity fracture in a double-porosity reservoir with a sealing fault boundary.

### 6.3. Available Theory and Solutions

### 6.3.1. Homogeneous reservoir

Based on Earlougher (1977) the classical equation which describes isothermal radial flow of a fluid throughout a homogeneous and isotropic medium, can be written as

$$
\begin{equation*}
\delta^{2} P / \delta r^{2}+1 / r \quad \delta P / \delta r=\left(\phi C_{t} \mu / k\right) \delta P / \delta t \tag{10}
\end{equation*}
$$

This expression is called the diffusivity equation, it assumes Darcian flow of a fluid of slight compressibility, through a medium of constant thickness, due the presence of a small pressure gradients. The term $\left(\mathrm{k} / \varnothing \mathrm{C}_{\mathrm{t}} \mu\right)$ is called hydraulic diffusivity. Solution of the diffusivity equation for the case of constant flow rate production in an infinite reservoir can be written as

$$
\begin{equation*}
P_{i}-P(r, t)=(q \mu / 4 \pi k H)\left[-E i\left(-\varnothing C_{t} \mu r^{2} / 4 k t\right)\right] \tag{11}
\end{equation*}
$$

where Ei is the exponential integral. When the exponential integral argument is lower than 0.01 , it can be approximated by

$$
\begin{equation*}
-\operatorname{Ei}\left[-\left(\varnothing C_{t} \mu r^{2} / 4 k t\right)\right] \approx \operatorname{Ln}\left[4 k t / \operatorname{Exp}(\Gamma) \varnothing C_{t} \mu r^{2}\right] \tag{12}
\end{equation*}
$$

Substituting Equation 12 into Equation 10 and remembering that $q=$ wv, gives

$$
\begin{equation*}
P_{i}-P(r, t)=(w v \mu / 4 \pi k H) \operatorname{Ln}\left[4 k t / \operatorname{Exp}(\Gamma) \varnothing C_{t} \mu r^{2}\right] \tag{13}
\end{equation*}
$$

for the case of a well ( $r=r_{w}$ ) which produces from all the reservoir thickness, and considering the skin factor
$P_{w f}=P_{i}-(w v \mu / 4 \pi k H)\left[\operatorname{Ln}\left(4 k t / \operatorname{Exp}(\Gamma) \varnothing C_{t} \mu r_{w}{ }^{2}\right)+2 s\right]$

Defining dimensionless time as

$$
\begin{equation*}
t_{D}=k t / \varnothing C_{t} \mu r_{w}{ }^{2} \tag{15}
\end{equation*}
$$

the dimensionless radius by

$$
\begin{equation*}
r_{D}=r / r_{w} \tag{16}
\end{equation*}
$$

and the dimensionless pressure as

$$
\begin{equation*}
P\left(r_{D}, t_{D}\right)=(2 \pi k H / w V \mu)\left[P_{i}-P(r, t)\right] \tag{17}
\end{equation*}
$$

For the dimensionless case, in which the mechanical skin factor is considered, one has from Equation 17 the following

$$
\begin{equation*}
P\left(1, t_{D}\right)+s=P\left(t_{D}\right)+s=(2 \pi k H / w v \mu)\left[P_{i}-P_{w f}\right] \tag{18}
\end{equation*}
$$

Substituting Equation 18 into Equation 14 and canceling similar terms

$$
\begin{equation*}
P\left(t_{D}\right)=0.5 \operatorname{Ln}\left[4 k t / \operatorname{Exp}(\Gamma) \varnothing C_{t} \mu r_{w}{ }^{2}\right] \tag{19}
\end{equation*}
$$

Equation 19 can be re-written using Equations 15 and 16 as

$$
\begin{equation*}
P\left(t_{D}\right)=0.5\left(\operatorname{Ln} t_{D}+0.8091\right) \tag{20}
\end{equation*}
$$

It is important to note that Equations 19 and 20 do not take into account the skin factor. Equations 14 and 20 are the most common solutions to the diffusivity equation for the production case in dimensional and dimensionless forms, respectively. For the case of the total drawdown-buildup period and using the Superposition Theorem, the resulting expression is written as

$$
\begin{equation*}
(2 \pi k H / w v \mu)\left[P_{i}-P_{w s}\right]=P_{D}\left(t_{D}+\delta t_{D}\right)-P_{D}\left(\delta t_{D}\right) \tag{21}
\end{equation*}
$$

Employing 20 and the dimensionless time definition

$$
\begin{equation*}
(2 \pi k H / w v \mu)\left[P_{i}-P_{w s}\right]=0.5 \operatorname{Ln}[(t+\delta t) / \delta t] \tag{22}
\end{equation*}
$$

Rearranging and changing the logarithmic base in Equation 22

$$
\begin{equation*}
P_{w s}=P_{i}-0.1832(w v \mu / k H) \log [(t+\delta t) / \delta t] \tag{23}
\end{equation*}
$$

On semilog paper, this equation describes a straight line with slope

$$
\begin{equation*}
m=0.1832(\omega v \mu / \mathrm{kH}) \tag{24}
\end{equation*}
$$

Now, substituting $P_{i}$ from Equation 18 into Equation 23 and solving for s

$$
\begin{gather*}
s=1.1513\left[\left(P_{w s}(t=1)-P_{w f}(8 t=0)\right) / m+\log \left(\left(t_{p}+1\right) / t_{p}\right)-\right. \\
\left.\log \left(k /\left(\phi C_{t} \mu r_{w}{ }^{2}\right)\right)-0.3513\right] \tag{25}
\end{gather*}
$$

Equations 23 and 25 are the basic expressions in the interpretation of buildup tests.

### 6.3.2. Distance to sealing fault

For pressure buildup testing, Earlougher (1977) used the intersection of the late time two straight lines and relates it to the dimensionless pressure, at the intersection time by

$$
\begin{equation*}
P_{D}\left(t_{D} /\left(2 L / r_{w}^{2}\right)\right)=0.5 \operatorname{Ln}\left[\left(t_{p}+\delta t\right) / \delta t\right]_{x} \tag{26}
\end{equation*}
$$

Thus, to estimate the distance to a linear fault, one finds [( $\left.\left.t_{p}+\delta t\right) / \delta t\right]$ when the semilog straight lines intersect and calculate $P_{D}$ from equation 26. From Figure 22 of the Earlougher monograph with that value of $P_{D}$ and the value of ( $t_{D} /\left(2 L / r_{w}{ }^{2}\right)$ is found, where finally

$$
\begin{equation*}
L=\sqrt{L}\left[k H t_{p} / 4 \varnothing C_{t} H \mu\left(t_{D} /\left(2 L / r_{w}^{2}\right)\right]\right. \tag{27}
\end{equation*}
$$

### 6.3.3. Double-porosity reservoir

Deruyck et al. (1982) present the theory for double-porosity behaviour which can be applied to both naturally fractured and multilayered reservoirs. The diffusivity equation for the fissured medium according to the terminology of Gringarten (1982) becomes

$$
\begin{equation*}
\mathrm{k}_{\mathrm{f}} / \mu \nabla^{2} \mathrm{P}_{\mathrm{f}}=\left(\varnothing \mathrm{VC}_{t}\right)_{f} \delta \mathrm{P}_{f} / \delta \mathrm{t}-\mathrm{q}^{*} \tag{28}
\end{equation*}
$$

The counter part equation for the matrix medium can be written

$$
\begin{equation*}
\mathrm{k}_{\mathrm{m}} / \mu \nabla^{2} \mathrm{P}_{\mathrm{m}}=\left(\varnothing \mathrm{VC}_{t}\right)_{\mathrm{m}} \delta \mathrm{P}_{\mathrm{m}} / \delta t+\mathrm{q}^{*} \tag{29}
\end{equation*}
$$

$V$ is the concentration of one medium (i.e. the ratio of that medium to the bulk volume); $q^{*}$ is the interporosity flow, namely, the flow from the matrix to the fissure. It is assumed in this system of equations, that the reservoir is of infinite lateral extent; with closed top and bottom boundaries; the fluid is slightly compressible; and flow is single phase and laminar; the gravitational forces are negligible; and the pressure gradients are small; the porosity of either medium is independent
of the pressure variations; and finally, the flow to the well occurs via the most permeable medium only, the least permeable medium just acting as a source. Here, the interporosity flow parameter is defined as

$$
\begin{equation*}
\lambda=\alpha r_{w}^{2} \mathrm{k}_{\mathrm{m}} / \mathrm{k}_{\mathrm{f}} \tag{30}
\end{equation*}
$$

and represents the facility of the fluid to flow from the matrix to the fissure. The ratio of the storativity of the fissure system to the storativity of the total fissure-matrix system is defined

$$
\begin{equation*}
\Omega=\left(\phi \mathrm{VC}_{t}\right)_{\mathrm{e}} /\left(\left(\phi \mathrm{VC}_{t}\right)_{\mathrm{f}}+\left(\phi \mathrm{VC}_{t}\right)_{\mathrm{m}}\right)=10^{-\delta \mathrm{P} / \mathrm{m}} \tag{31}
\end{equation*}
$$

where $\delta \mathrm{P}$ is the pressure increment between the two straight lines. For the case of constant flow rate production, the solution in the Laplace space for the fissured part is given by
$P_{f D L}=k_{0}\left[\sqrt{ }(s f(s)) \quad r_{D}\right] /\left(s \sqrt{ }(s f(s)) \quad k_{1}[\sqrt{ }(s f(s))]\right.$
where

$$
\begin{equation*}
f(s)=(\Omega(1-\Omega) s+\lambda) /((1-\Omega) s+\lambda) \tag{33}
\end{equation*}
$$

and $k_{0}$ and $k_{1}$ are the modified Bessel functions of second kind, of zero and first order respectively.

### 6.4. Estimation of Double-Porosity Parameters

### 6.4.1. Horner method

From the Horner plot Figure 16, the average reservoir pressure can be estimated, extrapolating the late time straight line, until it intersects the pressure axis. That occurs when $\log \left[\left(t_{p}+\delta t\right) / \delta t\right] \approx 0$, implies that $\delta t \geqslant t_{p}$. The late time straight line can be expressed as

$$
P(\delta t)=5.695-0.4445 \log \left[\left(t_{p}+\delta t\right) / \delta t\right]
$$

Therefore, the average reservoir pressure becomes

$$
P(\infty)=5.695
$$

One can estimate the average temperature in a similar form as for the pressure, however, as one can see in Table 2 or in Figure 10, the temperature stabilizes after 1 hour. This permits the assumption that the average reservoir temperature is $265{ }^{\circ} \mathrm{C}$, in that part of the field. In advance, it is pointed out that the interpretation of this test be based on the thermodynamic state defined by the average pressure and temperature of the reservoir. On this basis, one has from the compressed liquid water part of the steam tables (Keenan et al., 1978), the following data

$$
\begin{aligned}
& \mathrm{h}=1158[\mathrm{~kJ} / \mathrm{kg}] \\
& \mathrm{v}=1.2874 \mathrm{E}-03 \quad\left[\mathrm{~m}^{3} / \mathrm{kg}\right] \\
& \mu=1.0086 \mathrm{E}-04 \text { [Pa-s] }
\end{aligned}
$$

It is interesting to note that the specific stagnation enthalpy corresponding to this thermodynamic state is lower than that estimated by the James method.

Returning to the Horner plot (Figure 12), the left hand parallel straight line represents the total system, its slope can be calculated as

$$
m=(5.34-5.327) / \log (7.625 / 11.6)=-0.0713
$$

Therefore, the conductivity of the medium using Equation 24 is

$$
\begin{aligned}
\mathrm{kH} & =0.1832(16.19)(1.2874 \mathrm{E}-03)(1.0086 \mathrm{E}-04) /(0.0713) \\
& =5.4016 \mathrm{E}-12\left[\mathrm{~m}^{3}\right] \quad(18000[\mathrm{md}-\mathrm{ft}])
\end{aligned}
$$

The omega parameter can be evaluated using Equation 31 as

$$
\Omega=10-0.07156 / 0.0713=0.1
$$

As one can be noted from Equation 25 , to estimate the skin factor it is necessary to know the porosity, total compressibility and thickness of the reservoir, which are not available.

Nevertheless, to have some idea about this parameter, the following values were assumed

$$
\begin{aligned}
\emptyset & =0.1 \\
\mathrm{C}_{\mathrm{t}} & =1.865 \mathrm{E}-09\left[\mathrm{~Pa}^{-1}\right] \quad\left(1.2859 \mathrm{E}-05[\mathrm{psi}]^{-1}\right) \\
\mathrm{H} & =50[\mathrm{~m}] \quad(164[\mathrm{ft}])
\end{aligned}
$$

Therefore $P(t=1)$ can be calculated using the slope $m$ with

$$
\begin{aligned}
P(t=1) & =P_{1}=5.327-0.0713 \mathrm{Log}(190801 / 11.6) \\
P_{1} & =5.0262
\end{aligned}
$$

Substituting values into Equation 25

$$
\begin{aligned}
& s= 1.1513[(5.0262-5.135) / 0.0713+\log (190801 / 190800) \\
&-\log \left(5.4016 \mathrm{E}-12 /\left((0.1)(1.0086 \mathrm{E}-04)(1.865 \mathrm{E}-09)(0.108)^{2}(50)\right)\right. \\
&-0.3513] \\
& s=-5.26
\end{aligned}
$$

To estimate the distance to sealing fault

$$
P_{D}\left(t_{D} /\left(2 L / r_{w}^{2}\right)=0.5 \operatorname{Ln}(6)_{x} \approx 0.9\right.
$$

Coming into Figure 22 (Figure C. 2 in Earlougher, 1977)

$$
\mathrm{t}_{\mathrm{D}} /\left(2 \mathrm{~L} / \mathrm{r}_{\mathrm{w}}{ }^{2}\right)=2.4
$$

so that

```
L}=\sqrt{}{[(5.4016E-12)(190800) / (4 (0.1)(1.865E-09)(50)
        (1.0086E-04)(2.4))]
    \approx 338 [m] ( 1108 [ft])
```


### 6.4.2. Type-curve

In Figure 17, fitting the pressure buildup data, to the doubleporosity type curve of Bourdet et al. (1983), the match point is found

$$
\begin{aligned}
t_{D} / C_{D} & =110 \\
P_{D} & =24.3 \\
\delta P & =1 \\
\delta t & =3600 \quad[s] \quad(1 \quad[h])
\end{aligned}
$$

The following parameters were also determined

$$
\begin{aligned}
\left(c_{D} \operatorname{Exp}(2 s)\right)_{f} & =100 \\
\left(c_{D} \operatorname{Exp}(2 s)\right)_{f+m} & =10
\end{aligned}
$$

and

$$
\lambda \operatorname{Exp}(-2 \mathrm{~s})=1 \mathrm{E}-04
$$

Thus, the permeability-thickness (conductivity) of the most permeable system is

```
kf H = wv\mu/2\pi (PD
    =(16.19)(1.2874E-03)(1.0086E-04)(24.3)/ 2\pi(1E6)
        = 8.13028E-12 [m3 (27000 [md-ft])
```

The wellbore storage coefficient

$$
\begin{aligned}
c & =2 \pi k_{f} H / \mu\left(\delta t /\left(t_{D} / c_{D}\right)\right. \\
& =2 \pi(8.13028 \mathrm{E}-12)(3600) /(1.0086 \mathrm{E}-04)(110) \\
& =1.6756 \mathrm{E}-05\left[\mathrm{~m}^{3} / \mathrm{Pa}\right] \quad(0.7188 \text { barrel } / \mathrm{psi})
\end{aligned}
$$

the storativity ratio

$$
\begin{aligned}
\Omega & =\left(C_{D} \operatorname{Exp}(2 s)\right)_{f+m} /\left(C_{D} \operatorname{Exp}(2 s)\right)_{f} \\
& =10 / 100=0.1
\end{aligned}
$$

Using Equation 31, the storativity of the most permeable system is

$$
\begin{aligned}
\left(\emptyset \mathrm{VC}_{t}\right)_{\ell} & =\Omega\left(\emptyset \mathrm{VC}_{t}\right)_{t+\mathrm{m}} \\
& =0.1(9.325 \mathrm{E}-09)=9.325 \mathrm{E}-10[\mathrm{~Pa}]^{-1}
\end{aligned}
$$

The dimensionless wellbore storage coefficient can be calculated with

$$
\begin{aligned}
c_{D} & =c / 2 \pi r_{w}^{2}\left(\varnothing C_{t} H\right)_{t+m} \\
& =1.6576 \mathrm{E}-05 / 2 \pi(.108)^{2}(9.325 \mathrm{E}-09)=24255.125
\end{aligned}
$$

from

$$
\begin{aligned}
\left(C_{D} \operatorname{Exp}(2 s)\right)_{f+m} & =10 \\
s=0.5 \operatorname{Ln}\left(10 / c_{D}\right) & \approx-3.9
\end{aligned}
$$

Finally, the lambda parameter is

$$
\begin{aligned}
& \operatorname{Exp}(-2 \mathrm{~s})=1 \mathrm{E}-04 \\
& \lambda=1 \mathrm{E}-04 / \operatorname{Exp}(2(3.9))=4.0973 \mathrm{E}-08
\end{aligned}
$$

## 7. DRAWDOWN SIMULATION

The basic reservoir parameters from a pressure buildup test were determined using the available reservoir theory and thermodynamic principles. The methods of analysis applied to test data and the computed values of the reservoir parameters are summarized immediately

| Analysis Method | kH | $\Omega$ | $\mathrm{P}(\infty)$ |
| :---: | :---: | :---: | :---: |
|  | $\left[\mathrm{m}^{3}\right]$ X 1 E 12 | (dimensionless) | $[\mathrm{MPa}]$ |


| Horner Analysis | 5.4 | 0.1 | 5.7 |
| :--- | :---: | :---: | :---: |
| Type-Curve Analysis | 8.1 | 0.1 | - |

The skin factor, the length to the fault, and the interporosity flow coefficient, respectively, were estimated as

$$
\begin{aligned}
& \mathrm{s}=5.3 \quad \text { (Horner method) } \\
& \mathrm{s}=3.9 \quad \text { (type-curve) } \\
& \mathrm{L}=337.9 \quad \text { (Earlougher method) } \\
& \lambda=4.1 \mathrm{E}-08 \quad \text { (type-curve method) }
\end{aligned}
$$

These values were obtained by assuming

$$
\begin{aligned}
\varnothing & =10 \% \\
\mathrm{C}_{\mathrm{t}} & =1.9 \mathrm{E}-09 \quad\left[\mathrm{~Pa}^{-1}\right] \\
\mathrm{H} & =50[\mathrm{~m}]
\end{aligned}
$$

where $C_{t}$ is the liquid water compressibility (Grant et al., 1982) at average reservoir conditions. Thus the product ( $\varnothing \mathrm{C}_{\mathrm{t}} \mathrm{H}$ ) of the system, called total system storativity, was assumed as $9.3 \mathrm{E}-09 \mathrm{MPa}^{-1}$.

It is of interest to know and compare the ideal behaviour of the bottomhole pressure of a well in both homogeneous and double-porosity reservoirs, under similar conditions of
production to that of well A-18. For this purpose two different computer programs were made.

For the case of a well producing from a homogeneous reservoir, the exponential-integral solution was used as well as the superposition principle. The latter was done because of the necessity to take into account the sealing fault effect, which was substituted for an imaginary well producing at the same flow rate as the producing well.

In the case of a well producing from a double-porosity media, the equations given by Deruyck et al. (1982) and presented through this work, were included in the respective program. It was necessary to employ the Stehfest (1970) algorithm in order to find particular solutions to the fissured system. The superposition principle was also used to create the imaginary well to consider the linear boundary effect.

To be consistent in the drawdown simulation of the well under these two different conditions, the thermodynamic state of the fluid defined by the pressure and temperature at average reservoir conditions was used as well as the mass flow rate of well A-18 ( $16.2 \mathrm{~kg} / \mathrm{s}$ ). All the parameters obtained from the Horner method were utilized. The distance to sealing fault was that computed with the Earlougher method. To complete the data needed to simulate the well, the flow parameter coefficient $\lambda$ was that determined from type-curve analysis.

The program and the results of the simulation for the well in a homogeneous reservoir are given in Appendix B, while the program and results of simulation of the well in a doubleporosity reservoir are presented in Appendix C. The results for the homogeneous reservoir were plotted against the elapsed time and are shown in Figures 18 and 19. The results for the double-porosity reservoir were plotted against the elapsed time and are shown in Figures 20 and 21. The initial reservoir pressure was not plotted in these figures. Its value can be obtained by extrapolating the left hand straight
lines in Figures 18 and 20 to time zero. It was taken as 5.695 [MPa].

In Figures 18 and 20 (or Figures 19 and 21), it is possible to note at early time that the bottomhole pressure for the case of a well in a homogeneous reservoir decreases faster than the respective pressure for the case of a well in a double-porosity media. At later time, the bottomhole pressure declines smoothly in linear form.

To verify the responses of the simulated well in both reservoirs, it can be observed in Figures 19 and 21 that at early time the bottomhole pressure passes through a straight line with slope $m$. At the time when the boundary effects are felt, the pressure follows another straight with slope 2 m .

## 8. DISCUSSION

The theory presented above was developed considering idealized liquid phase flow. However, the existence of two-phase flow in the drawdown period was pointed out. When the superposition theorem was applied, a constant volumetric flow rate in both periods of the test (production-recovery) was assumed. The viscosity was also considered constant.

A double-porosity reservoir response and a sealing fault effects were detected in the pressure buildup test interpretation.

The simulation of an idealized well producing under similar conditions to that well A-18 from a double-porosity liquid reservoir was carried out for illustration purposes. However, the computed values of the bottomhole pressure for this mass flow rate ( $16.2 \mathrm{~kg} / \mathrm{s}$ ) can be considered valid within reasonable limits for a short production period. This assumption is based on a boiling front close to the well. Under shut-in conditions single liquid phase is found in the reservoir. For the specific production of well A-18 during the pressure test the induced boiling front was located inside the reservoir, as noted in the heat content increment in the output characteristic of Table 1. It is noted that the boiling front travels into the reservoir according to the extracted mass flow rate, so that for small flow rates as that produced for well A-18, the boiling front was located close to the well.

The simulation of this well for long time periods must be carried out using a more complex model (Grant et al., 1982) which includes relative permeability effects, two-phase compresibilities and rock heat transfer to the fluid.

## 9. CONCLUSIONS

1. Well A-18 produces from a single medium conductivity fracture located between 1200 and 1250 m depth at $265{ }^{\circ} \mathrm{C}$.
2. During the buildup test interpretation, it was possible to identify a double-porosity response
3. The late time data lie on a semilog straight line with a slope twice the total system slope. It was interpreted as a sealing fault boundary
4. Although the well produced from the reservoir twophases flow during the drawdown pressure period of the test, the single phase theory seems to be applicable for short time periods within reasonable limits
5. More complex models must applied to represent the phenomena of two-phase flow in the reservoir for both single-well and overall reservoir

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## NOMENCLATURE

```
A = head of water measured at the weir box [m]
C = compressibility [1/Pa]
c = wellbore storage [m3/Pa]
D = diameter [m]
H = reservoir thickness [m]
h = specific enthalpy [kJ/kg]
k = permeability [m
L = length to sealing fault [m]
m = slope
P = absolute pressure [MPa]
q = volumetric flow rate [m/s]
q}\mp@subsup{}{}{\mathbf{z}}=\mathrm{ interporosity flow [1/s]
r = radial distance [m]
s = skin factor (dimensionless)
T = temperature [ }\mp@subsup{}{}{\circ}\textrm{C}
t = time [s]
V = concentration of medium (dimensionless)
v = specific volume [m}\mp@subsup{\textrm{m}}{}{3}/\textrm{kg}
W = mass flow rate [kg/s]
X = steam fraction (dimensionless)
\alpha = geometrical factor (dimensionless)
\delta = increment or derivative or distance
\Gamma= Euler constant (0.57721)
\lambda = interporosity flow coefficient (dimensionless)
\mu = viscosity [Pa
\nabla = operator nabla
\Omega = ~ s t o r a t i v i t y ~ r a t i o ~ ( d i m e n s i o n l e s s )
\emptyset = ~ p o r o s i t y ~ ( d i m e n s i o n l e s s )
```


## SUBSCRIPTS

$0=$ stagnation
$1=$ taken at $t=1$
c = critical
b = dimensionless
d $=$ discharge
$f=$ most permeable media
i = initial
1 = liquid water
m = least permeable media
$\mathrm{p}=$ production
$s=$ steam
$t=$ total
$x=$ intersection
$w f=$ bottomhole (flowing)
ws $=$ bottomhole (static)
$w h=$ wellhead

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TABLE 1. Output curve data. (March, 1986)

| Pwh <br> $[\mathrm{MPa}]$ | STEAM <br> $[\mathrm{kg} / \mathrm{s}]$ | WATER <br> $[\mathrm{kg} / \mathrm{s}]$ | ENTHALPY <br> $[\mathrm{kJ} / \mathrm{kg}]$ |
| :---: | :---: | :---: | :---: |
| 0.9 | 25.6 | 16.6 | 1764 |
| 0.9 | 25.9 | 16.8 | 1763 |
| 1.9 | 25.3 | 17.5 | 1728 |
| 2.3 | 20.9 | 15.9 | 1676 |
| 3.2 | 6.6 | 9.6 | 1314 |

TABLE 2. Flowing pressure and temperature data (March 19, 1986)

| DEPTH <br> $[\mathrm{m}]$ | PRESSURE <br> [MPa] | TEMPERATURE <br> [C] |
| :---: | :---: | :---: |
| 0 | 3.236 | 237.88 |
| 100 | 3.349 | 241.26 |
| 200 | 3.468 | 244.47 |
| 300 | 3.591 | 246.77 |
| 400 | 3.709 | 248.31 |
| 500 | 3.827 | 249.99 |
| 600 | 3.958 | 251.83 |
| 700 | 4.085 | 253.21 |
| 800 | 4.222 | 254.90 |
| 900 | 4.358 | 256.43 |
| 1000 | 4.513 | 258.43 |
| 1050 | 4.615 | 259.50 |
| 1100 | 4.734 | 260.88 |
| 1150 | 4.849 | 262.13 |
| 1200 | 4.984 | 263.77 |
| 1250 | 5.282 | 262.46 |
| 1300 | 5.678 | 264.43 |
| 1320 | 5.822 | 266.57 |

TABLE 3. Pressure and temperature buildup (March 19, 1986)

| $\begin{array}{r} d t \\ {[\mathrm{~s}]} \end{array}$ | $\begin{array}{r} \mathrm{dt} \\ {[\mathrm{~h}]} \end{array}$ | PRESSURE [MPa] | TEMPERATURE <br> [C] | $\begin{gathered} \text { (tp+dt)/dt } \\ \text { (dimensionless) } \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0.000 | 5.135 | 263.12 |  | 0.000 |
| 360 | 0.100 | 5.271 | 264.27 | 531.000 | 0.136 |
| 480 | 0.133 | 5.282 | 264.43 | 398.000 | 0.147 |
| 600 | 0.167 | 5.295 | 264.43 | 319.000 | 0.160 |
| 720 | 0.200 | 5.299 |  | 266.000 | 0.164 |
| 840 | 0.233 | 5.303 |  | 228.143 | 0.168 |
| 960 | 0.267 | 5.307 |  | 199.750 | 0.172 |
| 1080 | 0.300 | 5.312 |  | 177.667 | 0.177 |
| 1200 | 0.333 | 5.316 |  | 160.000 | 0.181 |
| 1500 | 0.417 | 5.320 |  | 128.200 | 0.185 |
| 1800 | 0.500 | 5.323 |  | 107.000 | 0.188 |
| 2400 | 0.667 | 5.327 |  | 80.500 | 0.192 |
| 3000 | 0.833 | 5.332 |  | 64.600 | 0.197 |
| 3600 | 1.000 | 5.327 | 264.76 | 54.000 | 0.192 |
| 4800 | 1.333 | 5.323 |  | 40.750 | 0.188 |
| 6000 | 1.667 | 5.323 |  | 32.800 | 0.188 |
| 7200 | 2.000 | 5.320 | 264.76 | 27.500 | 0.185 |
| 8400 | 2.333 | 5.316 |  | 23.714 | 0.181 |
| 9600 | 2.667 | 5.316 |  | 20.875 | 0.181 |
| 10800 | 3.000 | 5.316 | 264.76 | 18.667 | 0.181 |
| 13200 | 3.667 | 5.320 |  | 15.455 | 0.185 |
| 15600 | 4.333 | 5.323 | 264.76 | 13.231 | 0.188 |
| 18000 | 5.000 | 5.327 | 264.76 | 11.600 | 0.192 |
| 21600 | 6.000 | 5.332 | 264.76 | 9.833 | 0.187 |
| 25200 | 7.000 | 5.336 |  | 8.571 | 0.201 |
| 28800 | 8.000 | 5.340 |  | 7.625 | 0.205 |
| 32400 | 9.000 | 5.344 | 264.76 | 6.889 | 0.209 |
| 36000 | 10.000 | 5.352 |  | 6.300 | 0.217 |
| 39600 | 11.000 | 5.361 | 264.76 | 5.818 | 0.226 |
| 43200 | 12.000 | 5.369 |  | 5.417 | 0.234 |
| 46800 | 13.000 | 5.381 | 264.76 | 5.077 | 0.246 |
| 50400 | 14.000 | 5.393 |  | 4.786 | 0.258 |
| 54000 | 15.000 | 5.406 | 264.76 | 4.533 | 0.271 |
| 61200 | 17.000 | 5.422 | 264.76 | 4.118 | 0.287 |
| 68400 | 19.000 | 5.438 | 264.76 | 3.789 | 0.303 |



FIGURE 1. MAIN GEOTHERHAL ZONES IN MEXICO


I IGURE ?. LOS AIUIRES CLOIIIFRMAL I IELD)


FIGURE 3. WELL COMPLETION


FIGURE 4. OUTPUT CHARACTERISTIC CURVES


FIGURE 5. FLOWING PRESSURE PROFILE


FIGURE 6. FLOWING TEMPERATURE PROFILE


FIGURE 7. THERMODYNAMIC STATES BEFORE SHUT-IN


FIGURE 8. COMPARISON BETWEEN MEASURED AND CALCULATED PRESSURE


Figure 9. Pressure buildup behavior


Figure 10. Temperature buildup behavior


FIGURE 11. THERMODYNAMIC BOTTOMHOLE STATE BEHACIOR DURING BUILDUP

Figure 12. Diagnostic plot


FIGURE 13. COMPARISON BETWEEN MEASURED AND FITTED PRESSURE INCREMENT


FIGURE 14. DIAGNOSTIC AND SPECIALIZED PLOTS

Figure 15. Specialized plot

$(t p+d t) / d t$
Figure 16. Horner. plot



Figure 18. Homogeneous reservoir

Figure 19. Homogeneous reservoir


Figure 20. Heterogeneous reservoir

Time [h]

FIGURE 22. Exponential-integral solution

Los Azufres well A-18

INPUT DATA AB FOLI.GH:

| WATER GRAUITY | 1.0000 |
| :--- | ---: |
| TOTSL HASS FLCSRATE,L.3/HR | 128516.0000 |
| HLAT TRANSF COEFF,BTU/HR/SQ | 0.0000 |

AT THE WELIMEAD :

| DEPTH,FT | 0.00 |
| :--- | ---: |
| PRESSURL,PSIA | 469.30 |
| TEHPERATURE,F | 460.62 |

PIPE DIAKETER USED AS FOLLCH:

$$
\begin{aligned}
& \text { FROM } \quad 0.0 \text { FT TO } 3146.0 \mathrm{FT}, \text { PIPE DJAMETER }(\mathrm{FT})=0.7296 \\
& \text { ARS ROUGHNESS }(\mathrm{FT})=0.0002 \\
& \text { FROK } 3145.0 \mathrm{FT} \text { TO } \begin{aligned}
4.331 .0 \mathrm{FT}, \text { PIPE DIAMETER }(\mathrm{FT}) & =0.5153 \\
\text { ABS ROUGHRESS }(\mathrm{FT}) & =0.0002 .
\end{aligned}
\end{aligned}
$$

TOIAL LENSTII DTUTDED IH 100 INTERUALS

DOUNHOLE SHUT-IN TEHPERATURE AS FOLLOH:

| DEPTH,FT | TFMF,F |
| :--- | ---: |
|  |  |
| 1540.00 | 482.00 |
| $229 \% .00$ | 489.01 |
| 2993.00 | 494.00 |
| 3445.00 | 499.00 |
| 3773.00 | 504.00 |
| 4101.00 | 504.00 |

1

| - TWO-PHASE FLOH * |  | FRICTION |  |  | ACCELE. Psi/100ft | POTENTIALPSi/100Pt | STM,FRAE: | REGIHE | $a w / A$$\mathrm{pt} / \mathrm{s}$ | $\begin{aligned} & \mathrm{as} / \mathrm{f} \\ & \mathrm{ft} / \mathrm{s} \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| DEPTH,FT | PRESIPSIA | TEMP, F | EN,BTU/L. 8 | isi/100ft |  |  |  |  |  |  |
| 0.00 | 469.30 | 460.62 | 565.00 |  |  |  | 0.1604 |  |  |  |
| 13. 31 | 171,18 | 451.03 | 365,00 | 0.2665 | 0.0000 | 4.0810 | 0.1599 | SLIJS | 1.3867 | 13.4831 |
| 86.62 | 473.07 | 461.44 | 565.00 | 0.2641 | 0.0000 | 4.0950 | $0.15 \% 4$ | SLIIG | 1.3890 | 13.3856 |
| 129,93 | 171.96 | 461,82 | 565,00 | 0.2618 | 0.0000 | 4.1095 | 0.1589 | Slug | 1.3892 | 13.2958 |
| 173.24 | 476,86 | 462.23 | 565,00 | 0.2594 | 0.0000 | 4.1242 | 0.1584 | SLUG | 1.3904 | 13.1995 |
| 216.35 | 478,77 | 4,52,64 | 585,00 | 0.2570 | 0.0000 | 4.1399 | 0,1579 | SLIJG | 1.3917 | 13.1039 |
| 259.86 | 480.68 | 463,05 | 565.00 | 0.2546 | 0.0000 | 4,1537 | 0.1574 | SLUG | 1.3930 | 13.0088 |
| 303.17 | 432.59 | 16.3 .47 | 565,00 | 0,2523 | 0,0000 | 4.1684 | 0.1568 | SLUG | 1.3943 | 12.9144 |
| 346.48 | 484.51 | 463.88 | 565.00 | 0.2499 | 0.0000 | 4.1831 | 0.1563 | SLIIE | 1.3956 | 12,8206 |
| 389.79 | 486,44 | 464,26 | 565.60 | 0,2176 | 0.0000 | 4.1979 | 0.1558 | SL.JS | 1.3969 | 12.7272 |
| 433.10 | 488.37 | 464.67 | 565.00 | 0.2459 | 0.0000 | 4.2198 | 0.1553 | Stug | 1.3981 | 12.6398 |
| 476.41 | 470,32 | 465,08 | 565,00 | 0.2143 | 0.0000 | 4.2176 | 0.1548 | GLJJG | 1.3994 | 12.5459 |
| 519.72 | 492.27 | 465.49 | 565.00 | 0.2428 | 0.0000 | 4.276 .3 | 0.1543 | SLUG | 1.4007 | 12,4523 |
| 563.03 | 491. 24 | 465.91 | 565,00 | 0.2112 | 0.0000 | 4,3047 | 0.1537 | SL.J: | 1.4020 | 12,3589 |
| 506.34 | 496.22 | 466.32 | 565.00 | 0.23\% | 0.0000 | 4.3337 | 0.1532 | SLIUG | 1.4033 | 12.2.658 |
| 619.65 | 493.21 | 466,74 | 565,00 | 0.2381 | 0,0000 | 4,3628 | 0.1527 | SLUG | 1.4046 | 12.1729 |


| 692.96 | 500.22 | 467.13 | 565.00 | 0.2366 | 0.0000 | 4.3903 | 0.1522 | SL.UG | 1.4050 | 12.0869 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 735.27 | 502,24 | 157.55 | 565,00 | 0.2.351 | 0.0000 | 4,1203 | 0.1516 | SLJJg | 1.4072 | 11.9946 |
| 779.58 | 504.26 | 467.97 | 565.00 | 0.2335 | 0.0000 | 4.1503 | 0.1511 | Slug | 1.4085 | 11.9025 |
| 822.89 | 506.30 | 453.39 | 565,00 | 0.2320 | 0.0000 | 4,4310 | 0.1506 | St, 106 | 1.4099 | 11.8106 |
| 866.20 | 508.36 | 468.81 | 565.00 | 0.2305 | 0.0000 | 4.5118 | 0.1500 | SLug | 1.4112 | 11.7190 |
| 909,51 | 510,43 | 469.24 | 565,60 | 0.2289 | 0.0000 | 4,5429 | 0.1195 | SLIJG | 1.4126 | 11,6277 |
| 952.82 | 512.51 | 469.67 | 565.00 | 0.2274 | 0.0000 | 4.574! | 0.1499 | SLUG | 1.4139 | 11.5366 |
| 996.13 | 514.60 | 470,06 | 565,00 | 0.2260 | 0.0000 | 4,6041 | 0.1484 | SLIJG | 1.4152 | 11.4522 |
| 1039,44 | 516.70 | 470.49 | 565.00 | 0.2245 | 0.0000 | 4.636, ${ }^{\text {a }}$ | 0.1479 | Stug | 1.4166 | 11.3616 |
| 1082.75 | 518.82 | 470,92 | 365.00 | 0.2230 | 0.0000 | 4.6689 | 0.1473 | SLug | 1.4180 | 11.2713 |
| 1126.06 | 520.95 | 471.35 | 565.00 | 0.2214 | 0.0000 | 4.7015 | 0.1467 | SLug | 1.4194 | 11.1812 |
| 1169.37 | 523.10 | 471,78 | 565.00 | 0.2199 | 0,0000 | 4.7349 | 0.1462 | SL.JI | 1.4208 | 11.0914 |
| 1212,68 | 525.26 | 472.22 | 565.00 | 0.2184 | 0.0000 | 4.7686 | 0.1450 | SL.UG | 1.4222 | 11.0018 |
| 1255,99 | 527,43 | 472,62 | 565.00 | 0.2169 | 0,0000 | 4,8027 | 0.1451 | SLUG | 1.4236 | 10.9125 |
| 1299.30 | 529.62 | 473.06 | 565.00 | 0.2155 | 0.0000 | 4.8316 | 0.1445 | SLUG | 1.4249 | 10.8295 |
| 1312.61 | 531,82 | 473,50 | 565.00 | 0.2141 | 0.0000 | 4,8695 | 0.1439 | SLIJS | 1.4264 | 10.7407 |
| 1385.92 | 534.04 | 473.94 | 565.00 | 0.212 .6 | 0.0000 | 4.9044 | 0.1434 | SL.JG | 1.4278 | 10.6522 |
| 1429,23 | 536.27 | 474,38 | 565.00 | 0.2111 | 0,0000 | 4.9401 | 0.1428 | SLIJS | 1.4292 | 10.5639 |
| 1472.54 | 538.51 | 474.82 | 565.00 | 0.2096 | 0.0000 | 4.975? | 0.1422 | SLUG | 1.4307 | 10.4758 |
| 1515.85 | 540.78 | 475,26 | 565.00 | 0.2032 | 0.0000 | 5.0125 | 0.1416 | SLJif | 1.4321 | 10.3880 |
| 1559.16 | 543.05 | 475.68 | 565.00 | 0.2068 | 0.0000 | 5.0467 | 0.1411 | SLUS | 1.4335 | 10.3063 |
| 1602.47 | 545,34 | 476,13 | 355.00 | 0.2053 | 0.0000 | 5,0341 | 0.1405 | SLIJS | 1.4350 | 10.2190 |
| 1645.78 | 547.65 | 476.58 | 565.00 | 0.2039 | 0.0000 | 5.1216 | 0.1399 | SLUG | 1.4365 | 10.1319 |
| 1689.09 | 549.97 | 477,03 | 565,00 | 0.202.4 | 0,0000 | 5.1599 | 0.1393 | SL.JG | 1.4380 | 10.0451 |
| 1732.40 | 552.31 | 477.48 | 565.00 | 0.2010 | 0.0000 | 5.1983 | 0.1387 | SLUG | 1.4395 | 9.9585 |
| 1775.71 | 551.66 | 477,73 | 565.00 | 0,1995 | 0.0000 | 5.2372 | 0,1381 | SI.US | 1.4410 | 9.8722 |
| 1819.02 | 557.03 | 478.39 | 565.00 | 0.1981 | 0.0000 | 5.2769 | 0.1375 | St.ug | 1.4425 | 9.7865 |
| 1362.33 | 559,42 | 473.81 | 565.00 | 0,1968 | 0,0000 | 5.3140 | 0.1369 | SUuG | 1.4439 | 9.7060 |
| 1905.64 | 561.83 | 479.27 | 565.00 | 0.1953 | 0.0000 | 5.3547 | 0.1363 | SLILS | 1.4454 | 9.6204 |
| 1943.95 | 561,25 | 179.73 | 555,00 | 0,1739 | 0.0000 | 5,3935 | 0.1357 | SLIUG | 1.4470 | 9.3350 |
| 1992.26 | 566.68 | 480.19 | 565.00 | 0.1925 | 0.0000 | 5.1367 | 0.1351 | Slug | 1.4485 | 9.4499 |
| 2035.57 | 559,14 | 480.66 | 565.00 | 0.1911 | 0.0000 | 5.1789 | 0.1344 | SL.UG | 1.4501 | 9.3651 |
| 2078.88 | 571.61 | 481.12 | 565.00 | 0.1887 | 0.0090 | 5.5212 | 0.1338 | SL.US | 1.4516 | 9.2805 |
| 2122.19 | 571.10 | 481,36 | 565,00 | 0.1832 | 0.0000 | 5.5641 | 0.1332 | slug | 1.4532 | 9.1961 |
| 2165.50 | 576.35 | 481.98 | 565.00 | 0.1670 | 0.0000 | 5.0078 | 0.1327 | SLUS | 1.4546 | 9.1219 |
| 2208.81 | 578,66 | 482,41 | 565.00 | 0.1703 | 0.0000 | 5,1793. | 0.1321 | SLIUS | 1.4561 | 9.0454 |
| 2252.12 | 581.06 | 482.85 | 565.20 | 0.1736 | 0.0000 | 5.354\% | 0.1315 | SLUIG | 1.4575 | 8.9666 |
| 2295, 43 | 583.53 | 483,31 | 565,00 | 0.1768 | 0.0000 | 5.5329 | 0.1308 | SLUS | 1.1591 | 8.38 .58 |
| 2338.74 | 586.08 | 483.78 | 565.00 | 0.1799 | 0.0000 | 5.7157 | 0.1302 | SL.Uls | 1.4607 | 8.8032 |
| 2382.05 | 588,72 | 481.27 | 585.00 | 0.1829 | 0.0000 | 5.7024 | 0.1295 | SLUG | 1.462 .3 | 8.7188 |
| 2425.36 | 591.44 | 484.74 | 565.00 | 0.1857 | 0.0000 | 6.073i | 0.1289 | Stug | 1.4640 | 8.6325 |
| 2469.67 | 374.24 | 185, 25 | 565,00 | 0.1883 | 0.0000 | 6,2759 | 0.1282 | SLIJG | 1.4656 | 8.5497 |
| 2511.98 | 597.12 | 485.77 | 565.00 | 0.1910 | 0.0000 | 6.4737 | 0.1275 | SLUG | 1.4674 | 8.4559 |
| 2555.29 | 600,10 | 486,31 | 565,00 | 0.1935 | 0.0000 | 6,6756 | 0,1267 | SLJg | 1.4692 | 8.3633 |
| 2598.60 | 603.16 | 486.87 | 565.00 | 0.1959 | 0.0000 | 6.8810 | 0.1260 | St.us | 1.4711 | 8.2750 |
| 2.341,91 | 606. 32 | 497.43 | 565,00 | 0.1981 | 0.0000 | 7.0900 | 0.1252 | scug | 1.4731 | 8.1800 |
| 2685.22 | 609.57 | 488.01 | 565.00 | 0.2002 | 0.0000 | 7.3024 | 0.1244 | SL.LIG | 1.4751 | 8.0832 |
| 2728.53 | 612,91 | 433.58 | 565,00 | 0.2020 | 0.0000 | 7.5071 | 0.1236 | SLIJ, | 1.4770 | 7.9900 |
| 2771.84 | 616.34 | 489.19 | 565.00 | 0.2037 | 0.0000 | 7.7260 | 0.1228 | St.ug | 1.4791 | 7.8901 |
| 2815,15 | 617,88 | 489,81 | 565,00 | 0.2053 | 0.0000 | 7.9481 | 0.1219 | Slug | 1.4813 | 7.7387 |
| 2858.46 | 623.50 | 490.45 | 565,00 | 0.2037 | 0.0000 | 8.1732 | 0.1210 | SLug | 1,4835 | 7.6858 |
| 2901.77 | 627,23 | 491,11 | 565,00 | 0.2079 | 0,0000 | 3. 1012 | 0,1201 | S1.JTS | 1,4358 | 7.5815 |
| 2945.08 | 631.06 | 491.74 | 565.00 | 0.2089 | 0.0000 | 8.6320 | 0.1192 | SLIIG | 1.4881 | 7,4757 |
| 2988.39 | 6314.99 | 472,42 | 565.00 | 0.2096 | 0,0060 | 8,8530 | 0.1182 | SL.UG | 1.4904 | 7.3734 |
| 3031.70 | 639.02 | 493.12 | 565.00 | 0.2103 | 0.0000 | 9.0907 | 0.1172 | SLUS | 1.4928 | 7,2652 |
| 3075.01 | 643,15 | 493.83 | 565,00 | 0.2107 | 0.0000 | 9.3289 | 0.1162 | SLUG | 1.4953 | 7.1357 |
| 3118.32 | 647,38 | 494.55 | 565.00 | 0.2109 | 0.0000 | 9.5692 | 0.1152 | SL.IJS | 1.4979 | 7.0451 |
| 31.51 .63 | 651,72 | 495.26 | 565,60 | 0.2109 | 0.0000 | 9.8116 | 0.1142 | SLUG | 1.5005 | 6.9334 |
| 3204.94 | 655.19 | 495.85 | 565.00 | 0.8903 | 0.0000 | 7.1072 | 0.1133 | SLug | 3.0127 | 13.7076 |
| 3248.25 | 658.68 | 476,44 | 565,00 | 0.8814 | 0,0000 | 7.1815 | 0.1125 | SLJJ | 3.0159 | 13.5308 |
| 3291.56 | 662.20 | 497.03 | 565.00 | 0.8726 | 0.0000 | 7.2567 | 0.1116 | SLUE | 3.0212 | 13.3518 |
| 3.331 .87 | 665.75 | 197.62 | 565.60 | 0.8337 | 0.0000 | 7,3332 | 0.1108 | SL.JG | 3.0255 | 13.1794 |
| 3378.18 | 669.33 | 498.22 | 565.00 | 0.8530 | 0.0000 | 7.4111 | 0.1099 | SLUE | 3.0299 | 13.0048 |
| 3421.49 | 672,94 | 498,32 | 565.60 | 0,8462 | 0.0000 | 7,1905 | 0.1091 | SLug | 3.0342 | 12,8309 |
| 3464.80 | 676.58 | 499,39 | 565.00 | 0.8380 | 0.0000 | 7,5670 | 0.1002 | slug | 3.0384 | 12,6663 |
| 3508.11 | 690.25 | 497,99 | 535,00 | 0.8293 | 0,0000 | 7.6491 | 0.1074 | SL.J3 | 3.0428 | 12.4937 |


| 3551.42 | 683.96 | 500.60 | 565,00 | 0.8206 | 0.0000 | 7,7329. | 0.1065 | SLIIS | 3.0473 | 12.3208 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3.591 .73 | 637,70 | 501.21 | 565.00 | 0.8119 | 0.0000 | 7.8188 | 0.1056 | Slus | 3.0518 | 12.1473 |
| 3638.04 | 691.47 | 501.82 | 565,00 | 0.8032 | 0.0000 | 7.9063 | 0.1017 | SLIJ | 3.0564 | 11.9745 |
| 3681.35 | 695,28 | 502.43 | 565.00 | 0.7946 | 0.0000 | 7,9955 | 0.1038 | SL.IJ | 3.0609 | 11.8021 |
| 3724.66 | 699.12 | 503.02 | 565.00 | 0.7880 | 0.0000 | 8.0863 | 0.1029 | SLUIG | 3.0556 | 11.6302 |
| 3767,97 | 703.00 | 303,64 | 565.60 | 0.7778 | 0,0000 | 8.1751 | 0.1020 | St.119 | 3.0700 | 11.4671 |
| 3811.28 | 706.91 | 504.27 | 565.00 | 0.7692 | 0.0000 | 8.2696 | 0.1010 | SLUG | 3.0747 | 11.2962 |
| 3854.59 | 710,86 | 501,90 | 565.00 | 0.7607 | 0.0000 | 8.3660 | 0.1001 | SL.JS | 3.0794 | 11.1258 |
| 3897,90 | 714.86 | 505.53 | 565.00 | 0.7522 | 0.0000 | 8.4644 | 0.0791 | slug | 3.0842 | 10.9558 |
| 3911.21 | 718,89 | 505.16 | 365.60 | 0,7137 | 0.0000 | 8.5653 | 0.0782 | SLJUG | 3.0890 | 10.7863 |
| 3984.52 | 727.96 | 506.80 | 565.00 | 0.7353 | 0.0000 | 8.6678 | 0.0572 | SLIGG | 3.0938 | 10.6172 |
| 4027.83 | 727,07 | 507,41 | 565,00 | 0.7268 | 0.0000 | 8.7726 | 0,0963 | SLUS | 3.0987 | 10.4485 |
| 4071.14 | 731.23 | 508.06 | 565.00 | 0.7188 | 0.0000 | 8.8745 | 0.0953 | SLUE | 3.1034 | 10.2882 |
| 4114.45 | 735,43 | 508.71 | 565.c0 | 0.7104 | 0.0000 | 8.9838 | 0,0943 | SLUE | 3.1084 | 10.1204 |
| 4157.76 | 739.67 | 507.36 | 565.00 | 0.7020 | 0.0000 | 9.0751 | 0,0933 | SL.UG | 3.1134 | 9.9530 |
| 4201.07 | 713,96 | 510,02 | 585,00 | 0,6937 | 0,0000 | 9,2095 | 0,0923 | SI,UG | 3.1185 | 9.7860 |
| 4244.38 | 748.30 | 510.68 | 565.00 | 0.6854 | 0.0000 | 9,32.66 | 0.0913 | SLUS | 3.1236 | 9.6194 |
| 4287.69 | 752,68 | 511.35 | 585,00 | 0.6771 | 0,0000 | 9,44.9 | 0.0903 | SLIIG | 3.1288 | 9.4531 |
| 4331.00 | 757.11 | 512,02 | 565.00 | 0.6638 | 0.0000 | 9.5681 | 0.0892 | SLU6 | 3.1340 | 9.2872 |

* PRESSURE AMALYSIS \&z

TOTAL FRICTIOA, LIQUID $=0.0000 \mathrm{FSI}$
TOIAL POTENTIAL, LIQUID $=0.0000$ PSI
TOTAL FRICTION, THO-PHACE $=15.8859$ PSI
TOTAL POIENTIAL, THO-PHASE $=271.9273$ PSI TOTAL ACCELE, TWO-PHASE $=0.0000$ FSS

```
INTEGER IOFT,IER
    realt8 madei, ard, y, x
iofl=2
```

THIS FRDGRAM IS USED TO IETEGHIHE THE DVEAL RCHAVIOR
OF A WELL FRODUCING AT CONSTAHT MpsS FLDU RAIE FFOII
A HOHOGENEDUS RESERVOIR WITH A SEAL ING FAULT BOUMHARY,
HERE 15 CALCULAIED THE EXF OSEHTIAL-INIEGKAL IO FIND IHE
FLOHING BOITOHHOLE FRESSURE.
Fo= Forosity (dimetsiorless)
$d v=$ dynanic viscosity (Fa-g)
tce lotal compreseesibility (1/fa)
co= conductivity (ntes)
$\mathrm{rt}=$ reservoir thickriess ( m )
fle nass flow rale (ks/s)
sve specific fluid volume (att3/ks)
wre wellbore radius ( m )
of $=$ distance to the sealing fault (b)
Fe= peraesbillity (matiz)
-
Fhi $=3.141592651$
$\mathrm{pi}=5.695$
$P 0=1$
$d v=1,0086 e-01$
lc=1,865e-09
$c 0=5,4016 \mathrm{e}-12$
rl. 50 .
$11=16.19$
$\mathrm{sv}=1.2874 \mathrm{e}-03$
wr=, 108
$d f=337.9$
$\mathrm{pe}=\mathrm{co} / \mathrm{rl}$
al=f1tsufdu/(1, ectafhitrefrl)
lype thal
be=pofdufle/pe
de=hefurti2./4.
el-belditiz.
do $20 \quad 1=1,10$
$t i=10.911$
ars $=\mathrm{de} / \mathrm{li}$
$y=$ andelllof harspier)
type tis
ars=et/ti
$x=$ andel(lopharsiler)
lyre fis
fufzpi-alf(yfy)
lype filipufivis
writello,'(4e15,6)') tifufigex.
20 continue
stop
end

```
Implicit realt8 (a-h),(c-z)
comaon/blkcf/v(20),miv
external pid
```

C this frogran can se usen to retermine the laealizen behavior
C OF A WELL Producing at constint hass flou rate from a
C HETEROGENEOUS (DOURLE-FOROSITY) RESERMOIR HITH
C A SEALINB FAULT BOURDARY.
C HERE IS USED THE STEHFEST ALGORITHM TQ FIND THE BDTTOHHOLF FRESSLIRE

## $n v=18$

1anbda $=4,0973 \mathrm{e}-08$
buega $=1$
wre. 108
$r \mathbf{r}=108$
df $=337.9$
$p i=5.695$
$11=16.19$
$5 \mathrm{v}=1.2874 \mathrm{x}-03$
$d v=1,0086 e-04$
phi $=3.141592654$
$\mathrm{co}=5.4016 \mathrm{e}-12$
al $=$ f1tsv*dv*1,e-06/(2,ifhi*(a)
call array (vinu)
do $i=1$ no
type tov(j), i
enddo
type tial
do $1 \quad 1=1,10$
do $2 \mathrm{j}=1$ 19
$t=10.2 x(i-1) \pm j$
pera/up
a=pfd(t,rolasbdaromeda)
$\mathrm{r}=2$, 1df/ur

pfi=pi-alt (atb)
Type totipli,arb
urite(10,'(4e15.6)') herfis3,h
2 continue
1 continue
stop
end

SUBROUTINE ARRAY(V,H)
c
c evaluates the array v(J) For i-terms in an agyatotic.
C SERIF. EXPAMSION DF PRODABILITY DENSITY FLNCTION.
C FOR SINGLE FRECISION 8-DIOIT ARITHMETIC THE OPTIHIU
c Valle fir $N$ IS abdut $N=10$, for double prectsion
C arithmetic the optimuh is arout $n=18$, n hust be eveh.
C based ow 'algorithy 368' by ho steifest in corannicat-
C IONS OF THE ACK, VOL 13 ,NO.1, JAN 1970.
c
IMPLICIT REAL. 8 ( $\mathrm{A}-\mathrm{H}, 0-2$ )
DIMENSIUN $\mathrm{G}(20), \mathrm{V}(20)$,H(20)
$G(1)=1$,
NH=H/2
DO $10 \mathrm{I}=2 \mathrm{~N}$
$10 \quad G(1)-G(I-1) * I$
C
$H(1)=2 . / G(\mathrm{NH}-1)$
DC $20 \mathrm{I}=2, \mathrm{NH}$
FI=I
IF (I.EQ.NH) GOTO 15

GOTO 20
15 $H(I)=F I * * N H \in(2 * I) /(G(I) * G(I-1))$
20 continue

C
D0 $30 \mathrm{I}=1, \mathrm{H}$
V( 1 ) $=0$ 。
$K_{1}=(I+1) / 2$
K2=1
If (K2.GT.NH) K2 2 = NH
Dif $40 \mathrm{~K}=\mathrm{K} \mathrm{L}, \mathrm{K} 2$
If (2KK-I, EQ.0) GOTO 36
IF ( I, EQ, K) GOTO 37

COTO 40
$36 \quad V(I)=J(I)+H(K) / G(I-K)$
COTO 40
$37 \quad V(\mathrm{I})=V(\mathrm{~J})+1 \mathrm{H}(\mathrm{K}) / 6(2 \times \mathrm{K}-\mathrm{I})$
4) continue
$y(\mathrm{I})=\mathrm{SH} \mathrm{HV}(\mathrm{I})$
SH:-5
C WKITE $(5,90)$ I, $\mathrm{V}(1)$
90 FСЕКАТ ( $3 \%, 16,3 \times, 613,6$ )
30 CONTINUE
RETURN
END

```
LPFD - The Laplace transfora of the pressure function for
                tis double porosity model.
    input parameters -
    S - The parzueler in the Laplace space.
    R - Radial distance,
    LAAMBDA -
    OHEGA -
    FPOGRAM SUBROUTINES - HMRSKO, RMRSK1, (The Bessel Punctions
        no and K1 frua the Ik,l library.
    FUNCTION LPFD(S, R, LAMBDA, OMEGA)
    IMP.ICIT REAL䇎 (A-H),(1)-Z)
    REAL&8 S, R, LAMBDA, OKEGA, LPFD, K0, K1, HAMBSKO, KHRSK1
    ENTERNAL MMBSKO, HHRSSK1
    FS = (OHEGA*(1, - OHECA)&S + LAKBDA) /((1,-OMEGA)&S + LAHBEA)
    X1 = SGRT(S*FS)
    x0 = x1*R
    ICPT = 1
    KO = MHBSKO(IOPT, XO, IER)
    K.1 = MmSSK1(TCPT, X1, tER)
    !PFD = K0/(S*X1KK1)
    RETHRN
END
```

```
`
C LPINV - Inverse Laplace transfuraction with the Stephest ilsoritha.
C
C INFUT PARAMETERS -
C . T-Time.
C R-Radial uistance.
C LAKBM -
c OMEGA -
C LPFD - The laplace transfora of the pressure function.
C OUTPUT PARAHETER -
C PFD - The pressure function.
SURROUTINE LPIN(T, R, LAHBDA, OKEGA, LPFD, PFD)
IMPLICIT REAL $8 (A-H), (O-Z)
REMLIS T, R, LAKBDA, OHEGA, LPFD, PFI
EXTERNAI. IPFD
COMHON/BLKCF/V(20), NV
xLN2 = D1.0s(2.00)
PP = XL.N2/T
PFD =0.
DO I = 1,NV
    S = PP&I
    PCD = PFI +V(I)&LPFD(S, R, LAKEDA, OHEGA)
END CO
PF!! = PFDPFP
RETUSN
EN%
```

C PFD - Dinensionless pressure for the double parosity acodel,
R - Radial distance,
LАĂDoA -
omega -

FUNCTION PFDCT, R, LAHBBM, DMEGA)

realzo $\mathrm{T}, \mathrm{R}$, LAMBda, OHEGA, PFn, LPFA EXTERNAL LPFD

## CALL LPINU(T, R, LAHBDA, DMEGA, LPFD, PFD)

RETURR
ENI

| .1000025+01 | $0.526693 E+01$ |  |  |
| :---: | :---: | :---: | :---: |
|  |  |  |  |
|  |  |  |  |
|  |  |  |  |
|  |  |  |  |
|  |  |  |  |
|  |  |  |  |
|  |  |  |  |
| 0.900000Etol |  |  |  |
| $0.100000 \mathrm{E}_{0} 0$ |  | 0.2 |  |
| 2000 | 0.51795 | 0.3 | .000000E+00 |
| 0.300000F.+02 | 0.5 |  |  |
| 0.400000Et02 | 0.515 |  |  |
| 0000Et02 | 0. |  |  |
| 0,600000E+02 | . 1.1 |  |  |
| 700000Et02 | 0.514106F+0 | 0.368306E+0 | 0.000000E+CO |
| 800000E +02 | 0.5136945 .01 | $0.374949 \mathrm{~F}+01$ | 0,000000E+00 |
| . $900000 \mathrm{E}+0$ ? | $0.513331 \mathrm{E}+0$ | 0.38 | 0.00000CE+00 |
| $0.100000 E+03$ | 0.513006E+01 | 0.386057E+0 | .000000Et+00 |
| 200000F.t03 | 0.51 | 0.4 |  |
| OODE +03 | 0.50 | 0.44 |  |
| .400000E+03 | 0,509722E+01 | $0.455219 \mathrm{E}+0$ | 0 |
| 0.500000E+03 | $0.508032 \mathrm{~F}+01$ | $0.466365 E+0$ | $0.7286768-37$ |
| DOOEF+03 | 0.50746 | 0. |  |
| OOOE +03 | 0.50599 | $0.4831755+$ |  |
| O000EF103 | 0.50S578E+01 | .489847Et |  |
| 0.900000e +03 | $0.506213 \mathrm{E}+0$ | 0.495733E+01 |  |
| 00 | 0.50 | . $500998 \mathrm{E}+$ | 0.203734E-27 |
| +04 | 0,503741E+0 | 0.535643E+0 | 0.102345E-20 |
| 300000E+04 | 0.502485E+ | 0.555 | .95811CE-18 |
| OOOOE+0 | 0.5015 | . 57 | . 5314 |
| 0.50000 E +04 | 0.5 | . $18814498+0$ |  |
| $0.600050 C+04$ | 0.5003 | 0.5 |  |
| 0.700000 E .104 | 0.4 | 0.598271 Et | $0.3115478-1$ |
| 0.8000)02 $5+04$ | $0.499440 E+0$ | $0.604947 \mathrm{E}+0$ | -0, 345 |
| 0.900000E + C4 | 0.499093E+0 | $0.610836 \mathrm{E}+0$ | -0.1 |
| $0.1000005+05$ | $0.498757 \mathrm{E}+1$ | $0.616104 E+0$ | -0,5755688-12 |
| $0.2000005+10.7$ | 0,496611 | 0.6 |  |
| 0.300000E+05 |  | $0.671032 \mathrm{~F}+0$ |  |
| 4000008+05 | 0.492461E+01 | 0.685416E+01 | 0.308 |
| 05 | 0.493773E+01 | 0.696573E+0 | 0.67 |
| $0.600000 E+0.7$ | 0.19 | 05 | -0,831189E |
| 0.700000E+05 | 0. | . $7133965+1$ | $-0.542535 E-08$ |
| 200000E+05 | 0.4923175+0 | 0.720073E+ | . 1 |
| 0.900000E+OJ | $0.491952 \mathrm{E}+01$ | 0.72 | 0.812304[-06 |
| 0,100000E+06 | $0.491626 E+0$ | 0.731 | 0.26884CE-05 |
|  |  |  |  |
| 0.300000 E +06 | 0.48 | 0.785 |  |
| $0.400000 E+06$ | 0.487 | 0.8005 |  |
|  | 0.4864 | 0 |  |
|  |  |  |  |
|  | 0. | 0,328525E+ |  |
|  | 0,4847 | 035 |  |
| $0.900000 E+06$ | 0.48123 | 0.341 |  |
| 0.100000E+07 | 0.483791E+01 | 0.846359E+01 | $0.1136615+00$ |
| 107 | 0.480574F+01 | $0.881016 \mathrm{E}^{\text {O }}$ O | . 28 |
| +07 | 0.478 | 0, |  |
| D000¢ +07 | $0.476916 \mathrm{E}+$ | 0.915 |  |
| OOSOOE+07 | 0.4 |  |  |
| , 6000005+07 | 0.47 | $0.9357478+0$ |  |
| 700000E+07 | $0.473745 E+01$ | 0.94365 4E+01 | 0.762571E+00 |
| 800000EF57 | 0.472969E+01 | $0.9503315+01$ | $0.821145 E+00$ |
| $0.900000 E+07$ | 0.472279E+01 | 0.956220E+01 | 0.8736177.+50 |
| 0.100009E+08 | 0,471658E+01 | $0.961488 E 101$ | 0,921150EF-O |
| Eto8 | 0.467511E+01 | 0.996145E+01 | $0.1244112+01$ |
| DOEF+08 | 0.465049E+0 | $0.101642 \mathrm{E}+$ |  |

